1.5 GPa compact double-wall clamp cell for SANS and NSE studies at low temperatures and high magnetic fields

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Abstract. We have developed and realised a compact high-pressure clamp cell of the piston-cylinder type for Small Angle Neutron Scattering (SANS) and Neutron Spin Echo (NSE) spectroscopy measurements. The cell can pressurise \emptyset 5 mm samples up to 1.5 GPa. It is non-magnetic, produces a sufficiently low and flat SANS background with a relatively good neutron transmission and can be used in combination with polarised neutrons. The cell accommodates samples of up to 55 mm³ in volume and adopts a two-layer geometry: a TiZr body (outer layer) and a CuBe2 insert (inner layer). Its design is very compact and can be employed inside cryomagnets featuring a sample bore greater than \emptyset 30 mm. We present the cell and illustrate its performances with a series of neutron scattering experiments performed on single-crystalline MnSi.

Keywords: High pressure, small angle neutron scattering, polarised neutrons

1. Introduction

Pressure is a fundamental thermodynamic variable and therefore plays an important role in condensed matter physics. It has a pronounced influence on the structural, magnetic, electric, chemical and mechanical properties of condensed materials [11]. In addition, the experimental investigation of some of the most fascinating phenomena in condensed matter physics, such as quantum critical points [3,6,15], hidden order [2,15] or non-Fermi liquid behavior [6], often require the combination of low temperatures, high magnetic fields and high pressures. In this field of research, neutron scattering is an invaluable probe as it provides direct microscopic information on both the structural and magnetic correlations. It is therefore important to develop non-magnetic pressure cells optimised for neutron scattering studies which can reach moderate pressures, beyond 1 GPa, whilst being able to accommodate relatively large sample volumes.

These requirements are met using piston-cylinder type clamp cells where the pressure is fixed by two knots. In addition, cells of the piston-cylinder type are compact enough to fit in the cryomagnets and cryostats typically used

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for neutron scattering experiments. In the following we present a new non-magnetic high pressure double-wall cell, which we have developed for Small Angle Neutron Scattering (SANS) and Neutron Spin Echo Spectroscopy (NSE) experiments. The cell has been successfully tested at the Institute Laue-Langevin (ILL) and the ISIS neutron source. It has been used to investigate the magnetic properties of the archetype helimagnet MnSi beyond the critical pressure of ~1.4 GPa where a non-Fermi Liquid behaviour sets-in [17]. We show that this pressure cell produces a low-enough and flat neutron background in the *Q*-range covered by typical SANS instruments and that it can also be used for NSE. The cell can accommodate samples up to 55 mm³ in volume and reaches pressures as high as 1.5 GPa at low temperatures.

2. Choice of materials

The choice of materials is the most important factor determining the performances of high pressure cells. In the specific case of neutron scattering, the materials used should not only have suitable mechanical properties including high tensile strength and hardness, but also be relatively transparent to hot, thermal and cold neutrons. In this way both the attenuation of the intense neutron beam and the activation of the pressure cell are minimised. Furthermore, it is important to choose materials with low coherent and incoherent scattering cross sections to reduce the scattering by the cell, and therefore decrease the (Q-dependent) background, or noise, of the measurement. To be able to apply magnetic fields or use polarised neutron beams (essential for NSE spectroscopy), it is of crucial importance that the chosen materials remain non-magnetic down to the lowest temperatures of interest.

Table 1 provides an overview of relevant material properties for chemical elements and the alloys considered for the construction of the pressure cell. The CuBe2 (2 wt.% Be), TiZr [22] and NiCrAl [19] alloys are all non-magnetic. Among them, the Russian Ni alloy with designation 40HNU-VI and mass weight composition provided in Table 2 has the highest hardness and has already been used for high-pressure neutron diffraction studies [4,19]. However, this alloy combines relatively high incoherent and coherent cross sections with strong forward scattering likely due to dislocations and/or short range order [18]. It is therefore much less appropriate for SANS than CuBe2 and TiZr.

The null-matrix TiZr alloy has a zero coherent cross-section σ_{coh} thanks to the combination of the negative and positive coherent scattering lengths b_{coh} of titanium and zirconium respectively (molar mass ratio 3.44/7.16,

		-				-	
Material	b _{coh} [fm]	$\sigma_{\rm coh}$ [barn]	$\sigma_{\rm inc}$ [barn]	$\sigma_{\rm total}$ [barn]	$\sigma_{\rm abs}$ [barn]	Hardness [HRC]	UTS [GPa]
Al	3.449	1.495	0.0082	1.503	0.231		
Be	7.79	7.63	0.0018	7.63	0.0076		
Cr	3.635	1.66	1.83	3.49	3.05		
Cu	7.718	7.485	0.55	8.03	3.78		
Ni	10.3	13.3	5.2	18.5	4.49		
Ti	-3.438	1.485	2.87	4.35	6.09		
Zr	7.16	6.44	0.02	6.46	0.185		
TiZr	0.06	0.04	314.8	314.9		32	0.86
CuBe2	7.73	7.50	0.49	7.99		35	1.20
NiCrAl	7.36	8.31	141.1	149.4		57-59	1.90

Table 1

Relevant properties for several materials considered for the pressure cell. Data on the scattering length and neutron cross section are adapted from ref. [21]. Hardness data are given in Rockwell units. Ultimate Tensile Stress (UTS) data at 300 K are provided in [16,24]

Table 2								
Relative mass weight composition of the Russian NiCrAl alloy 40HNU-VI								

С	Si	Mn	Fe	Ni	Cr	Мо	Ti	Al	Nb	Co	S	Р
≤0.03	≼0.1	≼0.1	≼0.6	balanced	39–41	-	-	3.3-3.8	-	-	≼0.01	≤0.01

i.e. 2.08 Ti/1 Zr). Unfortunately, this combination also leads to a relatively high incoherent cross-section σ_{inc} . The neutron scattering from TiZr is therefore almost *Q*-independent (flat) but high. Fully hardened CuBe2 has a comparable hardness but a much lower incoherent cross-section and therefore leads to a much lower *Q*-independent neutron background. The strong inelastic contribution observed at meV energy transfers will not be an issue for our needs [8]. Based on this information, we have decided to develop a two-layer clamp cell made from a CuBe2 insert pressed into a TiZr body.

3. Design of the cell

Figures 1, 2 and 3 show drawings and photographs of the double-wall cell and its auxiliary equipment. It consists of a hardened CuBe2 cylindrical insert pressed into a TiZr cylinder. The TiZr body is 54 mm long and has a maximum outer diameter of 29.5 mm. This diameter is reduced to 22.2 mm at the beam height on 5 mm. Flat faces are machined to ease its operation. M20 threads are provided at both ends to lock the NiCrAl pistons. With this design, the clamp cell matches most of the cryomagnets provided by neutron facilities up to fields of 10 Tesla.

The insert, made from hardened CuBe2, has inner and outer diameters of 4.9 and 10 mm respectively. Its length was chosen to match the dimensions of single crystals that are typically available. Before assembly, the inner cylinder of the TiZr body was machined with a conical shape inclined by 1° with the main axis. The outer diameter of the CuBe2 insert was also machined with the same conical shape. The insert was then pressed into the body by means of a press with a force of a few tons, until the insert is centred. The advantage of this two-layer geometry is that it minimises the total amount of material in the neutron beam whilst allowing to reach 1.5 GPa. Indeed, the maximal stress occurs near the bore and we reduce the attenuation by using TiZr in regions where the stress is lower [9].

The sample with a maximum diameter of 3.7 mm and a maximum height of 5 mm is encapsulated in a cylindrical Teflon (PTFE) can placed in the heart of the pressure cell. This Teflon can is topped up with a hydrostatic pressure transmitting liquid, typically Fluorinert [5]. At both ends of the can, the sealing is realised with two hardened CuBe2 rings with an outer diameter of 4.9 mm, only slightly greater than the diameter of the bore. The shape of the internal diameter of each ring is inclined by 45° so that the pressure slightly enlarges the diameter of the ring, ensuring its tightness. The optimum thickness of the rings is ≈ 0.4 mm.

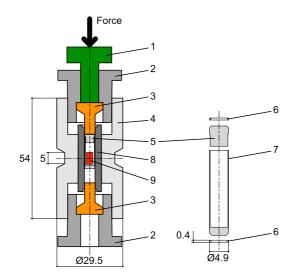


Fig. 1. Schematics of the assembled cylindrical multi-layer clamp cell: 1. NiCrAl piston, 2. hardened Ti nuts (grade VT8), 3. NiCrAl pistons, 4. TiZr body, 5. Teflon cap, 6. hardened CuBe2 rings, 7. Teflon can, 8. hardened CuBe2 insert, 9. sample. At the sample position, the diameter of the cell is reduced to 22.2 mm to improve the signal to background ratio.



(a) Side view

(b) Top view

Fig. 2. Photographs of the pressure cell, nuts and pistons.

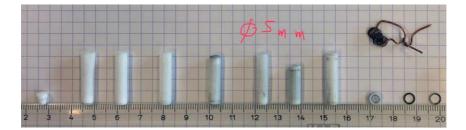


Fig. 3. Photographs of the PTFE sample cans and CuBe2 rings.

Nonmagnetic pistons made of NiCrAl alloy are positioned on both sides of the sample can. These pistons have a diameter which is 0.02 mm smaller than the inner diameter of the CuBe2 insert to avoid an excess of friction from the elastic deformation of the pistons at high pressures. Such a deformation ultimately results in "barrel-shaped" pistons leading to an increased contact area with the inner walls of the CuBe2 insert. This increase in friction lowers the effective pressure at the sample area and therefore reduces the effective operating pressure range. Molybdenum sulphide powder can be spread to reduce the friction with the inner surface of the CuBe2 insert but care must be taken because it activates in a neutron beam.

Finally, the pistons are clamped by two non-magnetic titanium (alloy grade VT8) knots positioned at both ends of the pressure cell. An M8 hole threaded in the top knot facilitates the fixation of the high pressure cell to the sample stick of the cryostat or cryomagnet.

4. Tests and operation of the cell

The cell has been used successfully for a series of neutron scattering experiments on the SANS instrument Larmor at ISIS (UK) and the NSE instrument IN15 at the ILL (France) [20]. The measurements were performed on a 110 mg single crystal of MnSi, approximately $2 \times 2 \times 3 \text{ mm}^3$ in size, originating from the same batch as the sample used in previous neutron scattering studies [1,13,14]. The sample was aligned with the [110] axis vertical, parallel to the Teflon sample can axis. This Teflon can was filled with chemically inert Fluorinert FC770 for the NSE experiments at IN15 and with a 50–50% mixture of FC770 and FC75 for the SANS experiments at Larmor [23]. Fluorinert has a very low incoherent scattering cross-section and thus does not contribute significantly to the

neutron background. This specific mixture was chosen because it features the highest hydrostatic limit [10]. During the filling of the can, care was taken to not encapsulate air bubbles. These bubbles collapse at low temperatures and significantly reduce the effective pressure on the sample.

At ISIS, we successfully applied pressures up to 2.2 GPa at room temperature. The pressure was applied using a 100 ton press and fixed by tightening the two titanium knots. These knots were tightened stepwise (i.e. first the top, then the bottom, then again the top, etc.) to prevent a displacement of the single crystal. The applied pressure was measured using a Strainstall 20 tons load cell. The press used at the ILL is equipped with an electronic load cell that accurately measures the load force of the press and the displacement sensor of the piston, which makes it possible to determine the friction and the pressure at room temperature by increasing and decreasing the load.

Figure 4 shows the relation between the applied force or the expected pressure in the cell at room temperature, and the effective pressure applied to the sample below 50 K. The pressure at low temperature is estimated from the critical temperature of MnSi as reported in ref [17], where the long range helimagnetic order sets-in at zero field. The figure shows that the pressure losses due to friction and the freezing of Fluorinert amount to about approximately 20%. These losses increase for higher applied pressures, possibly due to plastic instead of elastic deformation of the pistons.

Figure 5 presents the neutron transmission of the high pressure cell measured as a function of neutron wavelength at the time-of-flight instrument Larmor. The transmission decreases with increasing wavelength, from about 40% at $\lambda = 0.1$ nm to approximately 10% around $\lambda = 1.2$ nm. In addition, a dip is seen around $\lambda = 0.35$ nm, which is likely due to the Bragg scattering edge of the CuBe2 alloy. This is in good agreement with an estimated transmission $T = \exp[-\mu_{CuBe2}(10-4.9)] \times \exp[-\mu_{TiZr}(22.2-10)] = 12\%$ calculated with the linear absorption coefficients of CuBe2 and TiZr at $\lambda = 0.6$ nm: $\mu_{CuBe2} = 0.1563$ mm⁻¹ and $\mu_{TiZr} = 0.1020$ mm⁻¹.

The neutrons absorbed by the cell activate its constituents. We note that although Cu and Mn activate under neutron irradiation (both β -emitters), the activity outside the cell stayed below 2 μ Sv/h despite the long exposure time to the intense neutron beams of Larmor and IN15 (typically a week). The TiZr outer layer shields the activity originating from the decay of the irradiated Cu and Mn atoms. The resulting low activity measured outside the cell is a major advantage as it allows one to handle the cell and change the pressure directly immediately exposition to neutrons.

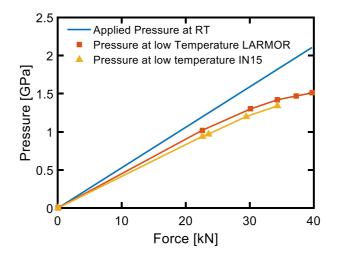


Fig. 4. Relation between the applied force or the expected pressure in the cell at room temperature, and the effective pressure applied to the sample below 50 K as deduced from the critical temperature of MnSi. The pressure at low temperatures is estimated from the critical temperature of the helimagnetic order at zero field as reported in [17].

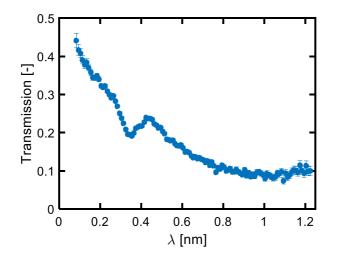


Fig. 5. Transmission of the empty cell as a function of the wavelength λ measured on Larmor using a 1 × 1 mm² neutron beam through the centre of the cell hosting a Teflon can without sample but with Fluorinert. The Teflon can being very thin (≈ 0.1 mm) and the Fluorinert FC75 and FC770 containing no hydrogen (generic composition (C_8F_{18})_n ($C_8F_{16}O$)_m), their contributions are quite low.

5. SANS and NSE experiments

The SANS measurements were performed on the time-of-flight instrument Larmor. The sample was illuminated by a $3 \times 3 \text{ mm}^2$ unpolarised incident neutron beam with a wavelength band of [0.09, 1.25] nm. The sample was positioned at 4.4 m from the detector consisting of 80 ³He position sensitive detector tubes, each 8 mm wide. The magnetic field was applied by a Scientific Magnetics 2T 3D vector field cryomagnet.

Figure 6 presents the scattering function S(Q) deduced by radially averaging the scattered intensity of the 2D SANS patterns, for the empty cryomagnet, the high pressure cell without sample and Fluorinert, and the high pressure cell with the MnSi sample at ambient pressure and 5 K. The curve obtained for the empty cryostat is typical of the background shape observed on all SANS instruments at ISIS, with a substantial background for $Q \leq 0.2 \text{ nm}^{-1}$. The plots show that the contribution of the cryomagnet to the background is negligible and that the background is mainly produced by the pressure cell. The low and flat background for $Q \gtrsim 0.2 \text{ nm}^{-1}$ allowed to get a high signal-to-noise ratio of 53.

To further illustrate the performance of the clamp cell, we present in Fig. 7 two SANS patterns of MnSi at 1.3 GPa and 8.3 K. Pattern (a) reveals two Bragg peaks indicative of the helimagnetic ordering at zero field. The SANS pattern (b) collected with a field 0.28 T applied along the incident neutron beam reveals six Bragg peaks, characteristic of skyrmion lattice correlations [7,12].

The NSE experiments were performed in the paramagnetic configuration on the NSE spectrometer IN15 using a 95% polarised incident neutron beam with a wavelength of $0.9 \pm 15\%$ nm. The NSE spectra were averaged over the entire detector. Figure 8 displays the intermediate scattering functions $I(\vec{Q}, t)$ as a function of the Fourier time t at 1.3 GPa (1.8 GPa at room temperature) for several temperatures below the critical temperature $T_c = 6.8$ K. As expected, $I(Q, t) \approx 1$ at all temperatures, which indicates that the helimagnetic correlations are completely elastic. These results confirm that the pressure cell does not depolarise the beam and can thus be used successfully in combination with polarised neutrons.

6. Conclusion

We have successfully designed, constructed and operated a compact high pressure cell of the piston-cylinder type with a relatively high transmission and sufficiently low background, adapted for neutron small angle scattering

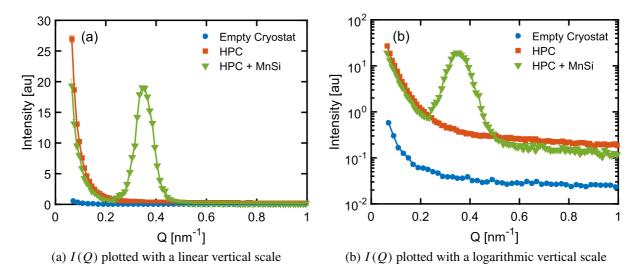


Fig. 6. Scattering function S(Q) measured on Larmor for the empty cryomagnet, the empty High Pressure Cell (HPC) and the MnSi sample at ambient pressure inside the HPC at 5 K and 0 T.

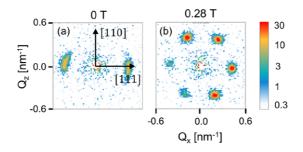


Fig. 7. Typical SANS patterns of MnSi at 1.3 GPa (1.7 GPa at room temperature) and T = 8.3 K. A magnetic field of (a) B = 0 T and (b) 0.28 T was applied along the incident neutron beam. The SANS patterns are background corrected using a high-temperature measurement at 50 K where the magnetic contributions to the scattering are negligible.

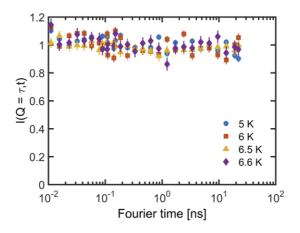


Fig. 8. Neutron spin echo spectroscopy results obtained on MnSi at a pressure of 1.3 GPa at zero magnetic field. The intermediate scattering function $I(\vec{Q}, t)$ is measured on the magnetic Bragg peak, i.e. $\vec{Q} = \vec{\tau}$, and is provided for four temperatures below the critical temperature of $T_c = 6.8$ K.

investigations at low temperatures and high magnetic fields. The high pressure cell is non-magnetic and can be used for investigations with polarised neutrons and NSE spectroscopy. The cell has a two-layer construction with a TiZr outer layer and a CuBe2 insert with an inner diameter of almost 5 mm. It can accommodate samples up to 55 mm³ in volume and is compatible with the cryomagnets typically used in neutron scattering. The high pressure cell can be pressurised up to 1.5 GPa and has been successfully used for SANS and NSE spectroscopy experiments on single-crystal MnSi.

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